

FIRST PRINCIPLES STUDY Ge -DOPED MONOLAYER GRAPHENE

S.S. HUSEYNOVA

*Institute of Physics, 131 H. Javid, Baku Az-1143, Azerbaijan**E mail: suma_huseynova.physics@mail.ru*

The electronic and structural properties of vacancy in graphene and doping of graphene with Ge calculated by the density functional theory (DFT) method within the Generalized Gradient Approximation (GGA). To simulate the vacancy effects have been studied 54 and 18 atom supercells to account for possible program software. Ferromagnetic spin ordering of vacancy of carbon atoms have been studied. It was shown that ferromagnetic ordering of the carbon atoms vacancies are located near the doping Ge atom. In addition DOS for 18 graphene atoms doped by Ge for s,p,d state have been calculated.

Keywords: *ab initio* calculation, DFT, ATK, graphene, Ge-doped, supercell, DOS, magnetic moment.

PACS: 31.10.+z, 31.15.E-, 75.50.Gg.

INTRODUCTION

Graphene is obvious to have important properties, in occasion of electronic conductivity, thermal stability, structural flexibility, and surface area that make it well suited as a building block in 3D architectures. The [1] result gives that the graphene sheets can improve the electrochemical performances of nanostructured transition metal oxides as anode materials for lithium-ion batteries. Graphene satisfies as conductive channels; graphene nanosheets facilitate charge transfer during the cycling process by forming an electrically conductive network with nanoparticles [2]. Graphene demonstrates high electronic performance [3–5] optical properties [6] thermal performance [7], larger specific surface area [8, 9] excellent elastic modules [10]. Graphene has an unusual band structure which is practically important for electronic devices. In [11] work all cases, along with lithium adsorption, the bandgap energy is increased, so that the germanium doped compound has the highest bandgap and the structure with no doped atom has the least bandgap. In [12] work have given for this material semiconducting properties and to rate its altering energy gap. Equivalent to other 2D structures, like germanene [13], silicene [14], and phosphorene [15], it is expected that graphene may be useful for the development of leading-edge technologies.

The experimental synthesis of germanium-graphene nanocomposite material [16], which can represent for lithium-ion battery applications [17], and the simulation of low energy ion implantation of germanium into a graphene target, which may be useful for single-atom catalysis [18]. Theoretical studies using *ab initio* calculations constitute a important and convenient method to obtain valuable information about nanomaterials [19–22]. Al-doped graphene was theoretically found to significantly increase the adsorbing energy of CO molecule, where CO binds to the top site of Al [23].

There are two types of graphene doping: electrical and chemical [24]. Have been investigated [25] the chemical doping of graphene, where doping occurs via chemical routes, namely substitutional doping with heteroatoms. In the present contribution, it is investigated the effects of the germanium doping

on the properties of graphene 54 and 18 atom supercells using first principles calculations.

MODEL AND CALCULATION METHOD

The calculations have been performed using the periodic using Atomistix Tool Kit (ATK) [26] implementing the spin-polarized DFT, PAW (projector-augmented wave method) and the PBE (Perdew–Burke–Ernzerhof exchange-correlation functional) [27].

Our calculations were performed for the primitive cell of Graphene and for a number of supercells with as many atoms as 54 by implementing the density functional theory (DFT) method within the Generalized Gradient Approximation (GGA) [28] and using the Atomistix Tool Kit program software. The Perdew–Burke–Ernzerhof (PBE) exchange-correlation functional (PBE) and Double Zeta Polarized basis sets were used in our calculations. The kinetic cut-off energy was 150 Ry. The primitive cell of Graphene was relaxed and optimized with force and stress tolerances of 0.01 eV/Å and 0.01 eV/Å³, respectively.

Ge DOPED AND VACANCIES IN GRAPHENE

To simulate the doping and vacancy effects we have studied bulk compound and its 54 atom graphene supercell. The following shares in the total magnetic moment are derived from Mulliken population analysis in the case of atom vacancy: 1.232 μ (μ_B), where μ_B is the Bohr magneton.

Fig. 1 shows the atomic structure of the vacancy in the 54 atoms graphene supercells and displays the density of states (DOS). The presence of carbon vacancies leads to acquired magnetic moment. 26th carbon atom acquires magnetic moment in the preferred direction $\mu(C) = (0.32 \mu_B)$. We established that acquired magnetic moments are: 3rd carbon atom ($0.32 \mu_B$), 21th ($0.32 \mu_B$), 26th ($0.327 \mu_B$), 9th ($0.101 \mu_B$), 15th ($0.101 \mu_B$) in the preferred direction, actually $\mu(C) = 8^{\text{th}}$ ($-0.035 \mu_B$), 31 ($-0.036 \mu_B$), 49nd ($-0.036 \mu_B$), and 14th carbon atom ($-0.033 \mu_B$) on the opposite direction.

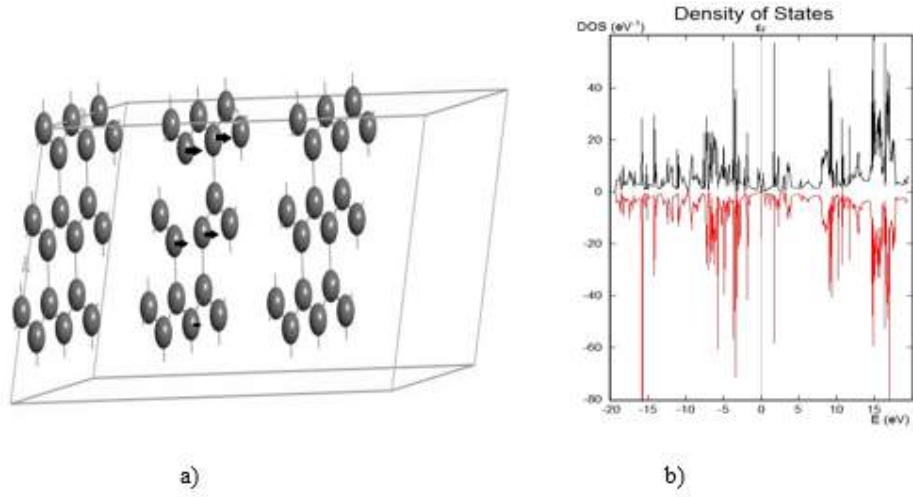


Fig. 1. The structure of defects in (a) supercell 54 atoms graphene with carbon vacancy and (b) DOS.

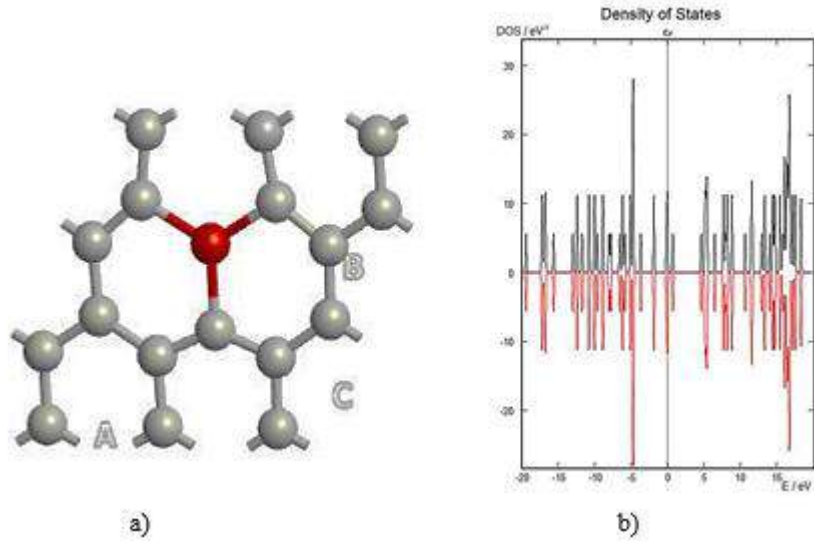


Fig. 2. (a) Atomic structure of 18 atom doped by Ge graphene and (b) density of state (DOS).

Fig. 2 displays the atomic structure of the Ge doped graphene plane and density of state (DOS).

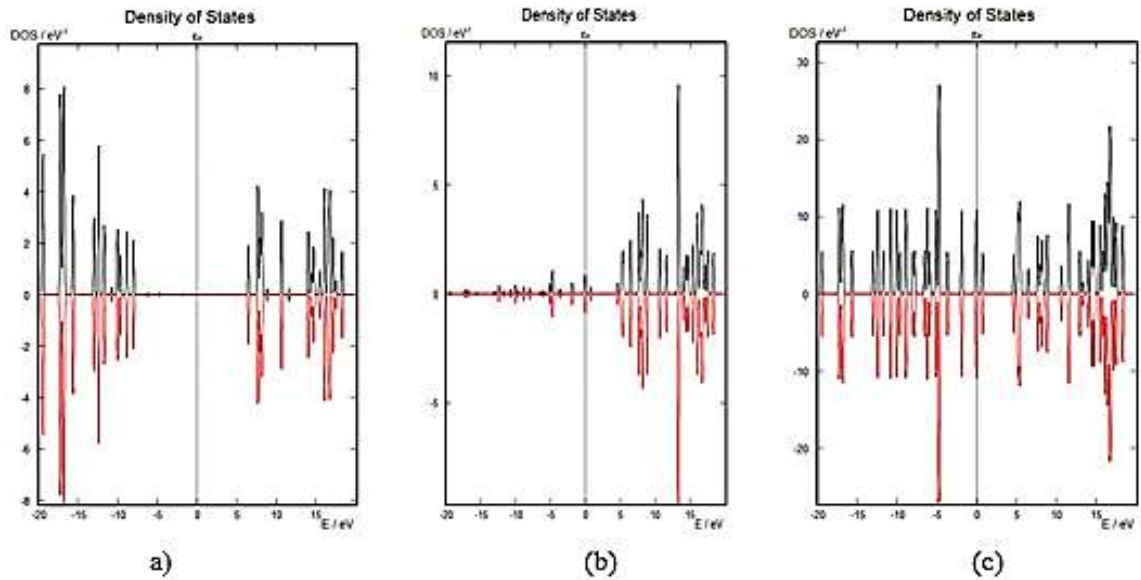


Fig. 3. The DOS for spin up and spin down d-, s-, and s p-electrons of the dopant Ge atom in the 18 atoms supercell: a) d- state, b) s-state, c) s p-state.

Fig. 3 shows the DOS for d-, s-, and -s p electrons of dopant Ge atom in the graphene supercells. The upper curve corresponds to electrons

with the direction of the spin up, and the lower with the direction of spin down.

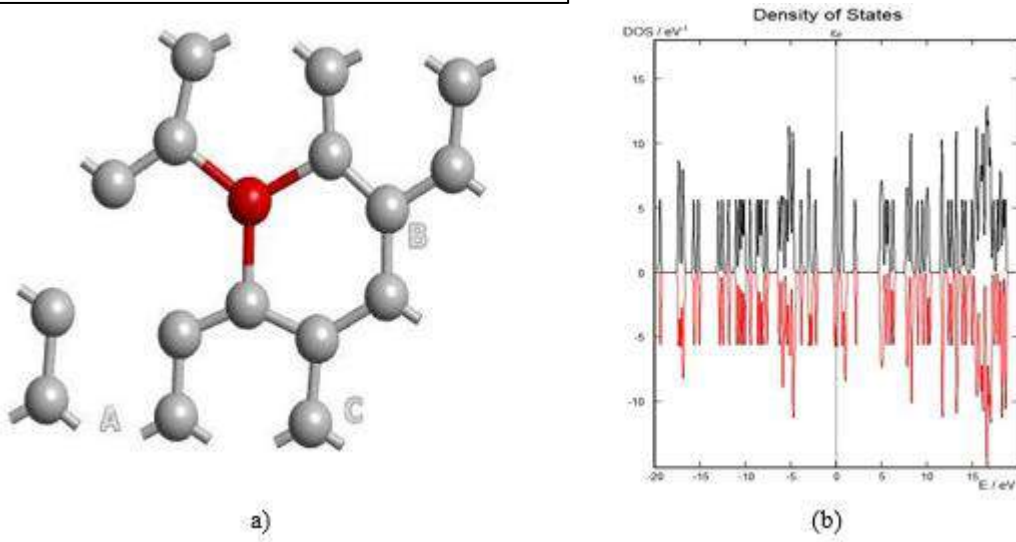


Fig. 4. (a) Atomic structure of the vacancy (Ge doped) and C vacancy in the graphene plane and (b) density of state (DOS).

Fig.4 displays the atomic structure of the Ge doped (a) graphene plane with C vacancy and (b) density of state (DOS).

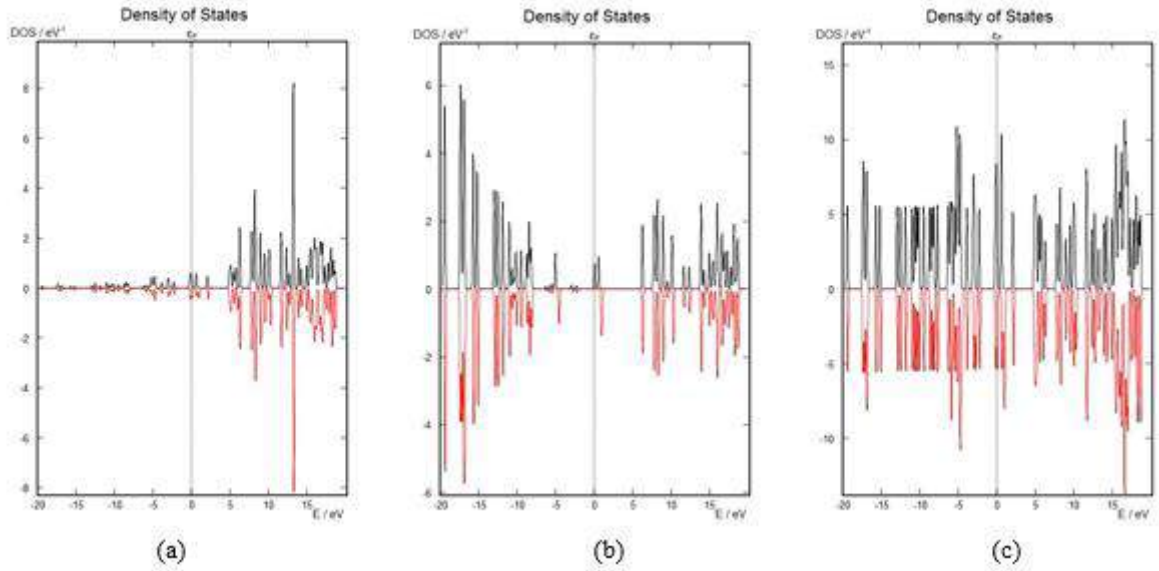


Fig. 5. The DOS for spin up and spin down d-, s-, and s p-electrons of the dopant Ge atom and C vacancy in the 18 atoms supercell : a) d-state, b) s-state, c) s p-state

Fig. 5 shows the DOS for d-, s-, and -s p electrons of dopant Ge atom and C vacancy in the graphene. Observed for the spin-down and spin up d-states below the Fermi level.

Si was the least effective element at opening the band gap, Ge ($r = 1.23 \text{ \AA}$) is a little more effective than Si ($r = 1.18 \text{ \AA}$), probably due to its larger size, which can induce more distortion than Si in the sp^2 C ($r = 0.77 \text{ \AA}$) framework [29].

The total magnetic moment in the supercell Ge doped graphene and C vacancy: $0.998 \mu (\mu_B)$. In case of Ge substitution C and in presence of C-vacancy which positioned near the dopant Ge atom creates magnetic moment $0.031 \mu_B$. Near the vacancy

$0.48 \mu (\mu_B)$ and the Ge additives $0.031 \mu (\mu_B)$, carbon atoms gained more magnetic moment.

CONCLUSIONS

The structural, electronic, and magnetic properties of Ge-doped graphene was studied using DFT calculations. Band structure and density of states Ge-doped, vacancy case of monolayer graphene are shown. To simulate the doping effect, we have performed our calculations for 54 atom graphene and Ge doped 18 atom-contained supercells. It has been established that doping the monolayer of graphene leads to a significant altering its electronic properties.

Mulliken population analysis shows that the presence Ge- doped and carbon vacancy state graphene supercell acquire $\mu(C) = (0.998 \mu B)$ magnetic moment.

The presence vacancy graphene supercell acquired magnetic moment of the carbon near the

vacancy is noticeably larger $\mu(C) = (0.998 \mu B)$ than that of the other without vacancy state $\mu(C) = (0.001 \mu B)$. In the presence of Ge-doped and no vacancies in supercell, $E_g = 0.505 \text{ eV}$, while Ge-doped and vacancies $E_g = 0.21 \text{ eV}$.

- [1] P. Lian, X. Zhu, H. Xiang, Z. Li, W. Yang, H. Wang. Enhanced cycling performance of Fe₃O₄ – graphene nanocomposite as an anode material for lithium-ion batteries. *Electrochim Acta*, 56:834-840, 2010. doi:10.1016/j.electacta.2010.09.086.
- [2] AK Rai, J. Gim, LT Anh, J Kim. Partially reduced Co₃O₄/graphene nanocomposite as an anode material for secondary lithium ion battery. *Electrochim Acta*, 100:63-71, 2013. A simple and rapid urea-assisted, auto-combustion approach for the synthesis of high-performance Co₃O₄/graphene nanocomposite anode is reported, DOI:10.1016/j.electacta.2013.03.140.
- [3] C.W.J. Beenakker. Colloquium: Andreev reflection and Klein tunneling in graphene, *Rev. Mod. Phys.*, 2008, 80, 1337–1354, DOI: <https://doi.org/10.1103/RevModPhys.80.1337>.
- [4] A.H. Castro Neto, F. Guinea, N.M.R. Peres, K.S. Novoselov and A.K. Geim. The electronic properties of graphene, *Rev. Mod. Phys.*, 2009, 81,109–162, DOI: 10.1103/RevModPhys.81.109.
- [5] S.D. Sarma, S. Adam, E.H. Hwang and E. Rossi. Electronic transport in two dimensional graphene, *Rev. Mod. Phys.*, 83, 407–470, 2011. DOI: <https://doi.org/10.1103/RevModPhys.83.407>.
- [6] R.R. Nair, P. Blake, A.N. Grigorenko, K.S. Novoselov, T.J. Booth, T. Stauber, N. MR. Peres and A.K. Geim. Fine structure constant defines visual transparency of graphene, *Science*, 320, 1308, 2008. doi: 10.1126/science.1156965.
- [7] A.A. Balandin, S. Ghosh, W. Bao, I. Calizo, D. Teweldebrhan, F. Miao and C.N. Lau. Superior thermal conductivity of single-layer graphene, *Nano Lett.*, 8, 902–907, 2008. doi/10.1021/nl0731872.
- [8] B. Luo, B. Wang, X.L. Li, Y.Y. Jia, M.H. Liang and L.J. Zhi. Graphene-confined Sn nanosheets with enhanced lithium storage capability, *Adv. Mater.*, 24, 3538–3543, 2012. DOI: 10.1002/adma.201201173.
- [9] C.H. Xu, B.H. Xu, Y. Gu, Z.G. Xiong, J. Sun and X.S. Zhao. Graphene-based electrodes for electrochemical energy storage, *Energy Environ. Sci.*, 6, 1388–1414, 2013 DOI: 10.1002/adma.201201173.
- [10] C. Lee, X.D. Wei, J.W. Kysar and J. Hone. Measurement of the elastic properties and intrinsic strength of monolayer graphene, *Science*, 321, 385–388. 2008. doi: 10.1126/science.1157996.
- [11] M.M. Loghavia, b, H. Mohammadi-Maneshb, R. Egraa, A. Ghasemia and M. Babaieea. DFT Study of Adsorption of Lithium on Si, Ge-doped Divacancy Defected Graphene as Anode Material of Li-ion Battery *Phys. Chem. Res.*, Vol. 6, No. 4, 871-878, December 2018. DOI: 10.22036/pcr.2018.148943.1543.
- [12] S.S. Huseynova. AB initio study of Si -doped monolayer graphene, *Transactions of National Academy of Science of Azerbaijan, series of physics-mathematical and technical, Physics and Astronomy*, № 2, 2020.
- [13] A. Acun, L. Zhang, P. Bampoulis, M. Farmanbar, A. van Houselt, A.N. Rudenko, M. Lingenfelder, G. Brocks, B. Poelsema, M.I. Katsnelson, and H.J.W. Zandvliet Germanene. The germanium analogue of graphene, *J. Phys.: Condens. Matter* 27, 443002 (11 pp), 2015. doi:10.1088/0953-8984/27/44/443002.
- [14] B. Lalmi, H. Oughaddou, H. Enriquez, A. Kara, S. Vizzini, B. Ealet, B. Aufray. Epitaxial growth of a silicene sheet, *Appl. Phys. Lett.* 97, 223109, 2010. <https://doi.org/10.1063/1.3524215>.
- [15] Z. Guan, W. Wang, J. Huang, X. Wu, Q. Li, J. Yang. Metal-Free Magnetism and Half-Metallicity of Carbon Nitride Nanotubes: A First-Principles Study, *J. Phys. Chem. C* 118, 39, 22491-22498, 2014. <https://doi.org/10.1021/jp508617k>.
- [16] E. Aktürk, C. Ataca, S. Ciraci. Effects of silicon and germanium adsorbed on graphene, *Appl. Phys. Lett.* 96, 123112, 2010; <https://doi.org/10.1063/1.3368704>
- [17] J.G. Ren, Q.H. Wu, H. Tang, G. Hong, W. Zhang, S.T. Lee. Germanium–graphene composite anode for high-energy lithium batteries with long cycle life, *J. Mater. Chem. A* 1, 1821 2013. <https://doi.org/10.1039/C2TA01286C>.
- [18] M. Tripathi, A. Markevich, R. Böttger, S. Facsko, E. Besley, J. Kotakoski, T. Susi. Implanting Germanium into Graphene *ACS Nano* 12, 4641, 2018. <https://doi.org/10.1021/acsnano.8b01191>.
- [19] M. Sajjad, N. Singh, U. Schwingenschlgl. Strongly bound excitons in monolayer PtS₂ and PtSe₂ *Appl. Phys. Lett.* 112, 043101 2018, <https://doi.org/10.1063/1.5010881>.

- [20] *N. Singh, Udo Schwingenschlögl.* A Route to Permanent Valley Polarization in Monolayer MoS₂, *Adv. Mater.* 29(1), 1600970, 2017. DOI: 10.1002/adma.201600970.
- [21] *M.H. Mohammed, A.S. Al-Asadi, F.H. Hanoon.* Superlattices Microstructures 129, 14 -19 (2019) Electronic structure and band gap engineering of bilayer graphene nanoflakes in the presence of nitrogen, boron and boron nitride impurities, DOI:10.1016/j.spmi.2019.03.012.
- [22] *M.H. Mohammed, A.S. Al-Asadi, F.H. Hanoon.* Solid State Commun. 282, 28-32, 2018. Semi-metallic bilayer MS₂ (M= W, Mo) induced by Boron, Carbon, and Nitrogen impurities, Solid State Communications, Pub.Date : 2018. DOI: 10.1016/j.ssc.2018.07.011.
- [23] *Z.M. Ao, J. Yang, S. Li and Q. Jiang.* Enhancement of CO detection in Al doped graphene *Chem. Phys. Lett.*, 2008,461, 276, <https://doi.org/10.1016/j.cplett.2008.07.039>.
- [24] *Lv Ruitao and M. Terrones.* Towards new graphene materials: Doped graphene sheets and nanoribbons, *Mater.Lett.*78, 209-218, 2012. DOI:10.1016/j.matlet.2012.04.033.
- [25] *M Yu Arsent'ev, A V Prikhodko, A V Shmigel, T L Egorova and M V Kalinina.* Doping graphene with a monovacancy: bonding and magnetism, *Journal of Physics: Conference Series* 661, 012028, 2015. doi: 10.1088/1742-6596/661/1/012028 <http://quantumwise.com>.
- [26] *M. Ernzerhofa and G.E. Scuseria.* Assessment of the Perdew–Burke–Ernzerhof exchange–correlation functional, *Journal of Chemical Physics*, v.110, No 11. 5029-5036, 1999.
- [27] *S.A. Tolba, K.M. Gameel, B.A. Ali, H.A. Almossalami and N.K. Allam.* The DFT+U: Approaches, Accuracy, and Applications, 2018. DOI:10.5772/intechopen.72020. <https://www.sciencedirect.com/topics/chemistry/generalized-gradient-approximation>, 1-30.
- [28] *A. Pablo Denis.* Chemical Reactivity and Band-Gap Opening of Graphene Doped with Gallium, Germanium, Arsenic, and Selenium Atoms, DOI: 10.1002/cphc.201402608, 2014. Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim.

Received: 10.03.2020